

# Photovoltaik 3

Electrons and Holes in Semiconductors

Resources:

Peter Würfel, Physik der Solarzellen

Jenny Nelson, The Physics of Solar Cells

A. Goetzberger, Sonnenenergy: Photovoltaik

Thomas Dittrich, lecture notes

1. Free electron and hole densities
2. Intrinsic carrier concentration
3. Doping of semiconductors
4. Equilibrium and excess charge carriers



non-occupied states

energy ↑

energy interval

$\Delta E$

$E_C$

conduction band

$E_g$

$E_V$

valence band

energy interval

$\Delta E$

occupied states



# Density of free electrons and holes

$$dn, dp \propto dE$$

$$dn \propto D_e(E)$$

$$dp \propto D_h(E)$$

$$dn \propto f_e(E) = f(E)$$

$$dp \propto f_h(E) = 1 - f(E)$$

E energy

dE energy interval

$D_{e,h}$  density of electron states

f occupation probability for an electron

$E_A$  width of the conduction and valence bands

$$dn = D_e(E) \cdot f(E) \cdot dE$$

$$n = \int_{E_C}^{E_C + E_A} D_e(E) \cdot f(E) \cdot dE$$

$$dp = D_h(E) \cdot [1 - f(E)] \cdot dE$$

$$p = \int_{E_V - E_A}^{E_V} D_h(E) \cdot [1 - f(E)] \cdot dE$$

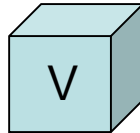
# Density of states

number of states per volume and energy interval

$$D(E) \equiv \frac{1}{V} \cdot \frac{dN(E)}{dE}$$

# Delocalization

Free charge carriers do not belong to a chemical bond.  
Free charge carriers in an ideal semiconductor are delocalized.  
One free charge carrier is delocalized over the whole volume.



$$(\Delta x)^3 = V$$

$\Delta x$  – volume of one delocalized state  
One delocalized state occupies the whole volume.

# Delocalization and Heisenberg uncertainty principle

Heisenberg uncertainty principle  
 $\Delta p$  momentum of an electron in a state  
 $h = 6.626 \cdot 10^{-34}$  Js

$$\Delta x \cdot \Delta p = h$$

phase space volume per state

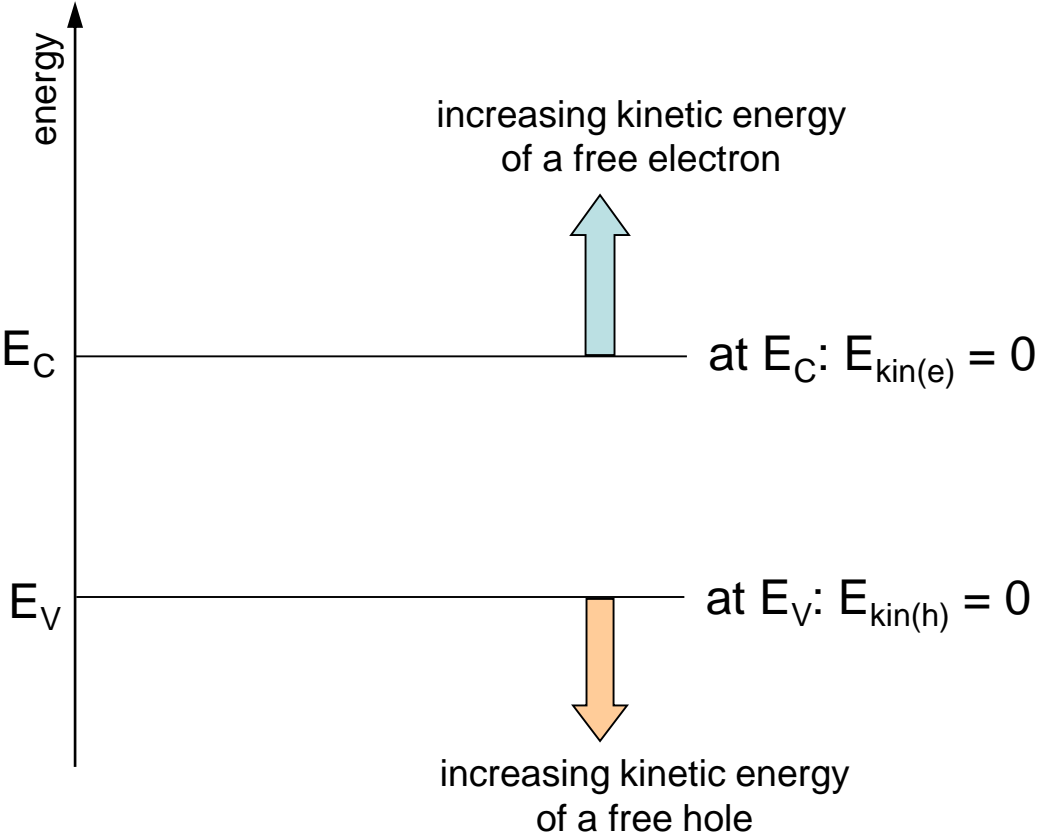
$$(\Delta x)^3 \cdot (\Delta p)^3 = h^3$$

volume in momentum space which  
can be occupied by one state

$$(\Delta p)^3 = \frac{h^3}{V}$$

keep in mind: p means here momentum

# Kinetic energy of free electrons and holes



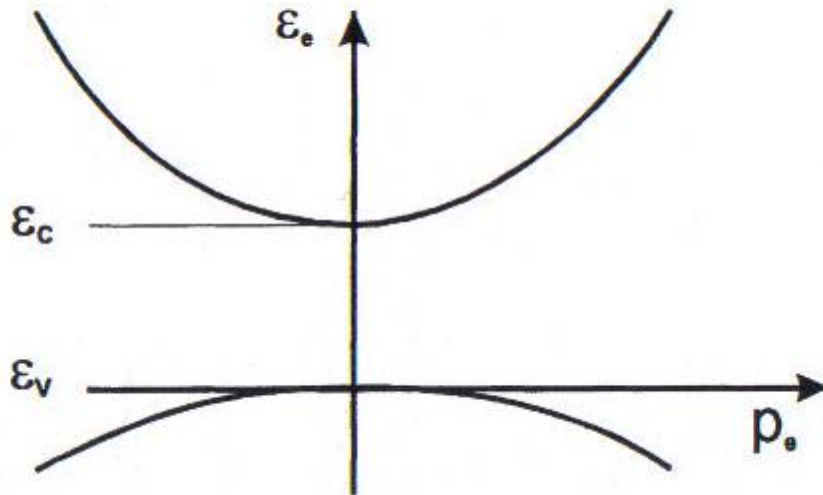
$$E_{kin} = \frac{p^2}{2 \cdot m}$$

p momentum  
m mass

keep in mind: p means here momentum



# Effective mass



$$\frac{1}{m_e^*} = \frac{\partial^2 \epsilon_e}{\partial p_e^2}$$

# Degeneration of states



non-degenerated state



threefold degenerated state

tenfold degenerated state



states in a cupboard: one place can be occupied by more than one cup or glass

degeneration: one state can be occupied by more than one particle

twofold degeneration of electron states due to spin-up and spin-down

# From the number of states as a function of momentum to the density of states as a function of energy

for states with  $p' \leq p$

$$N(|p|) \propto \frac{4\pi}{3} \cdot |p|^3 \quad \text{volume of a sphere with radius } |p| \text{ in momentum space}$$

$$N(|p|) \propto 2 \quad \text{spin degeneration}$$

$$N(|p|) \propto \frac{1}{(\Delta p)^3} \quad \text{inverse volume of one state in momentum space}$$

$$N(|p|) = \frac{8\pi}{3 \cdot h^3} \cdot V \cdot |p|^3 \quad \text{substitution: } p_{(e)} = \sqrt{(E - E_C) \cdot 2 \cdot m_e^*} \quad N_e(E) = \frac{8\pi}{3 \cdot h^3} \cdot V \cdot (2 \cdot m_e^*)^{3/2} \cdot (E - E_C)^{3/2} \quad D_e(E) = \dots$$

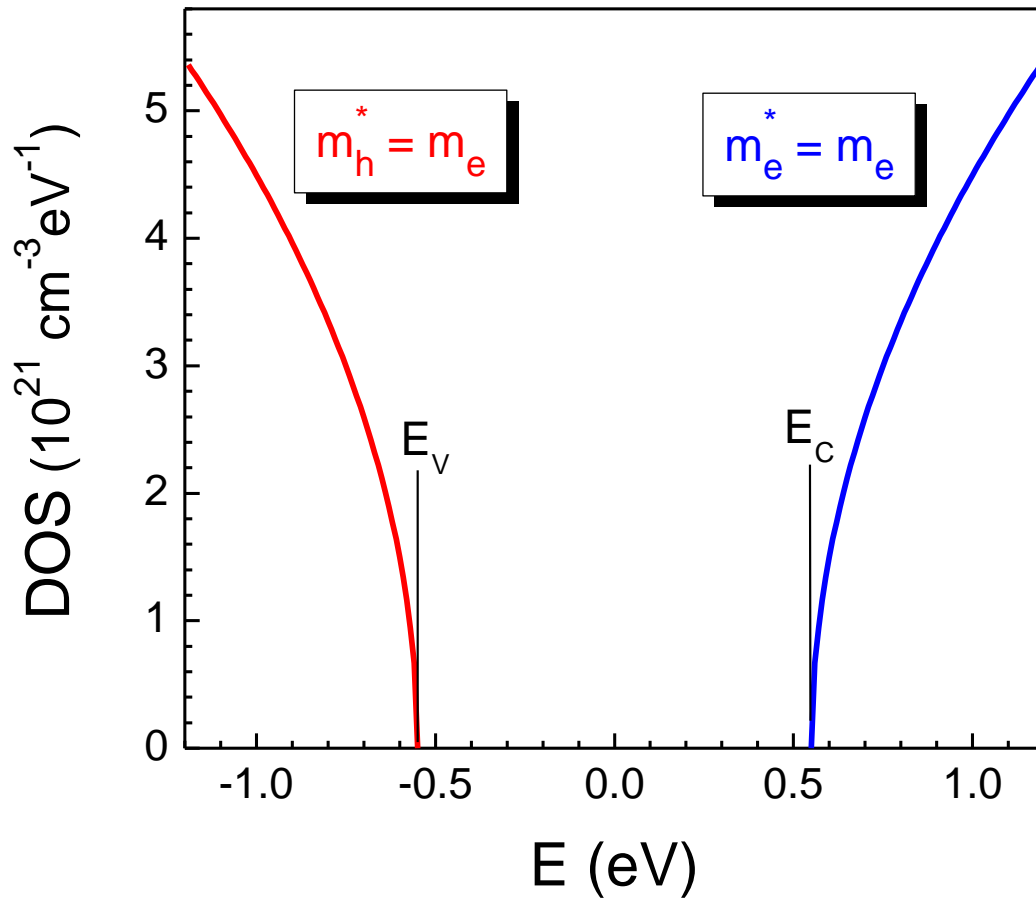
$$D_e(E) = 4\pi \cdot \left( \frac{2 \cdot m_e^*}{h^2} \right)^{3/2} \cdot (E - E_C)^{1/2}$$

$$D_h(E) = 4\pi \cdot \left( \frac{2 \cdot m_h^*}{h^2} \right)^{3/2} \cdot (E_V - E)^{1/2}$$

\*

keep in mind: p means here momentum

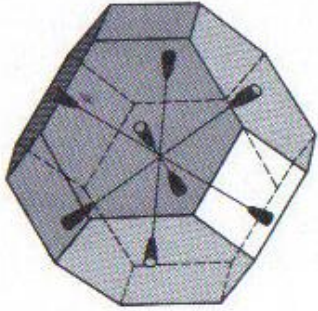
# Order of DOS in conduction and valence bands



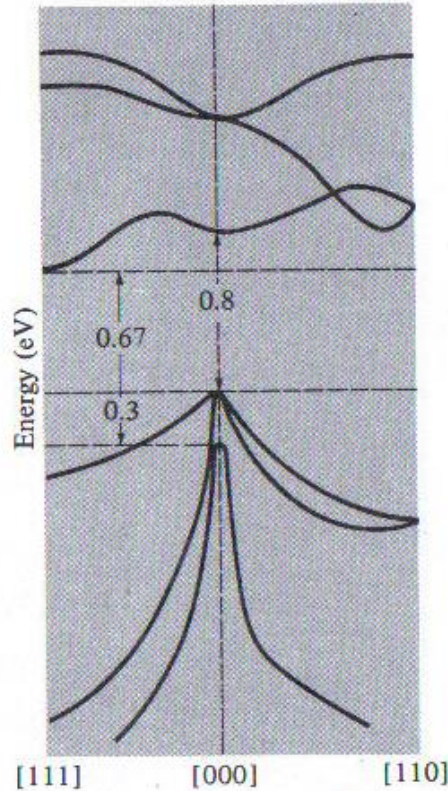
	$\frac{m_e^*}{m_e}$	$\frac{m_h^*}{m_e}$
Si	1.08	0.55
Ge	0.88	0.29
GaAs	0.067	0.47

$$D(E) = \frac{6.7 \cdot 10^{21}}{\text{cm}^3 \cdot \text{eV}} \cdot \left( \frac{m_e^*}{m_e} \right)^{3/2} \cdot (E - E_c)^{1/2}$$

# Band structure of Germanium



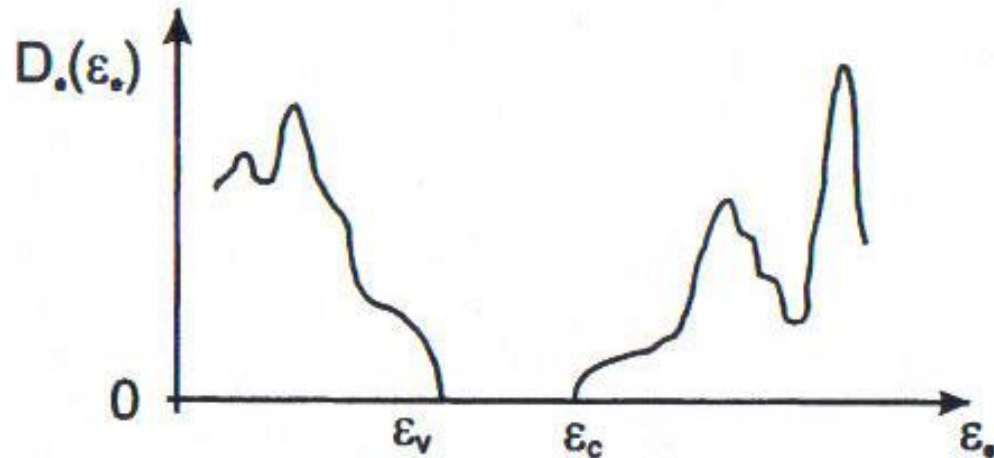
- Constant-energy surface near the conduction band minimum
- Representation by four ellipsoids



Energy bands of Ge along [111]

- Heavy and light holes
- For large k deviations from parabolic approximation

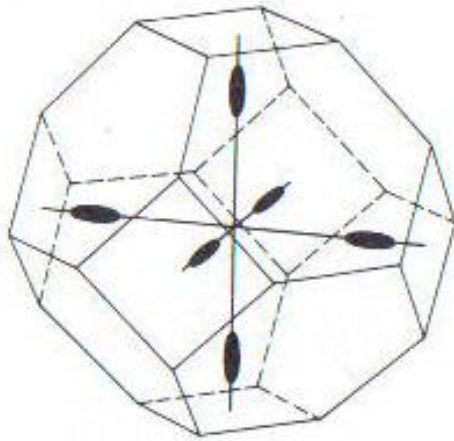
# Density of States of Germanium



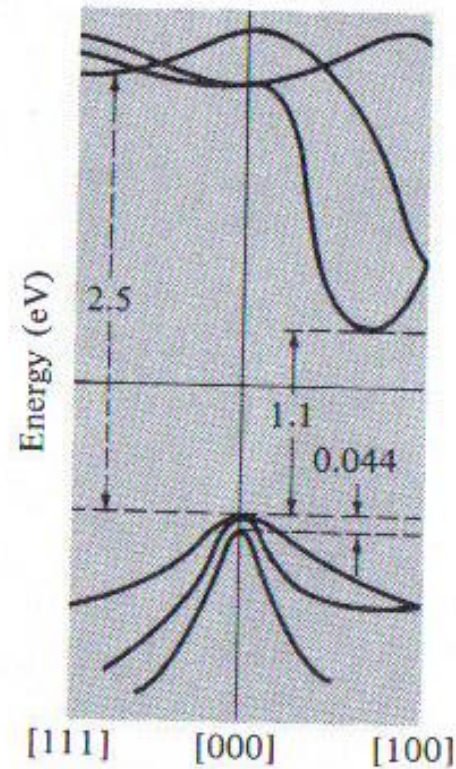
Zustandsdichte für Elektronen im Leitungs- und Valenzband des Halbleiters Germanium

# Band structure of Silicon

## Indirect semiconductor



Constant energy surface  
near conduction band minimum



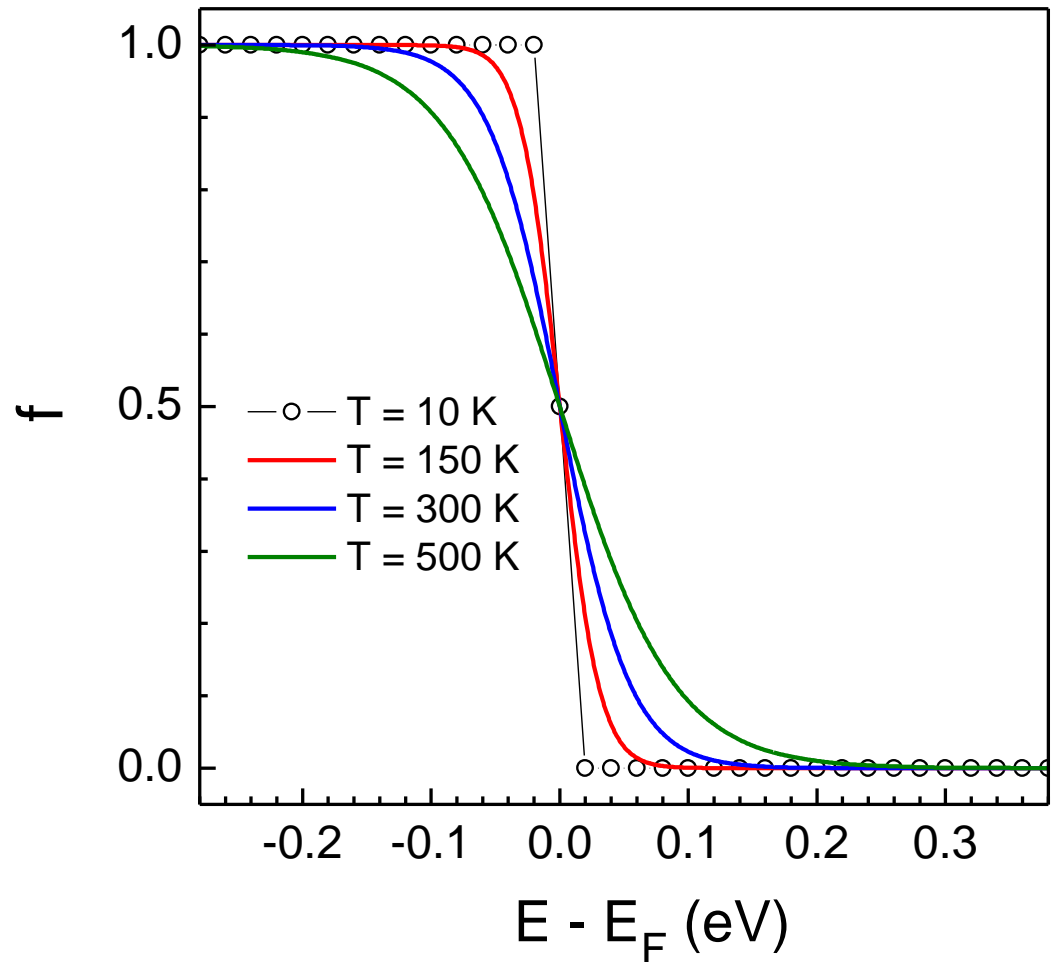
Energy bands in silicon

# Fermi-Dirac statistics

$$f(E) = \frac{1}{\exp\left(\frac{E - E_F}{k_B \cdot T}\right) + 1}$$

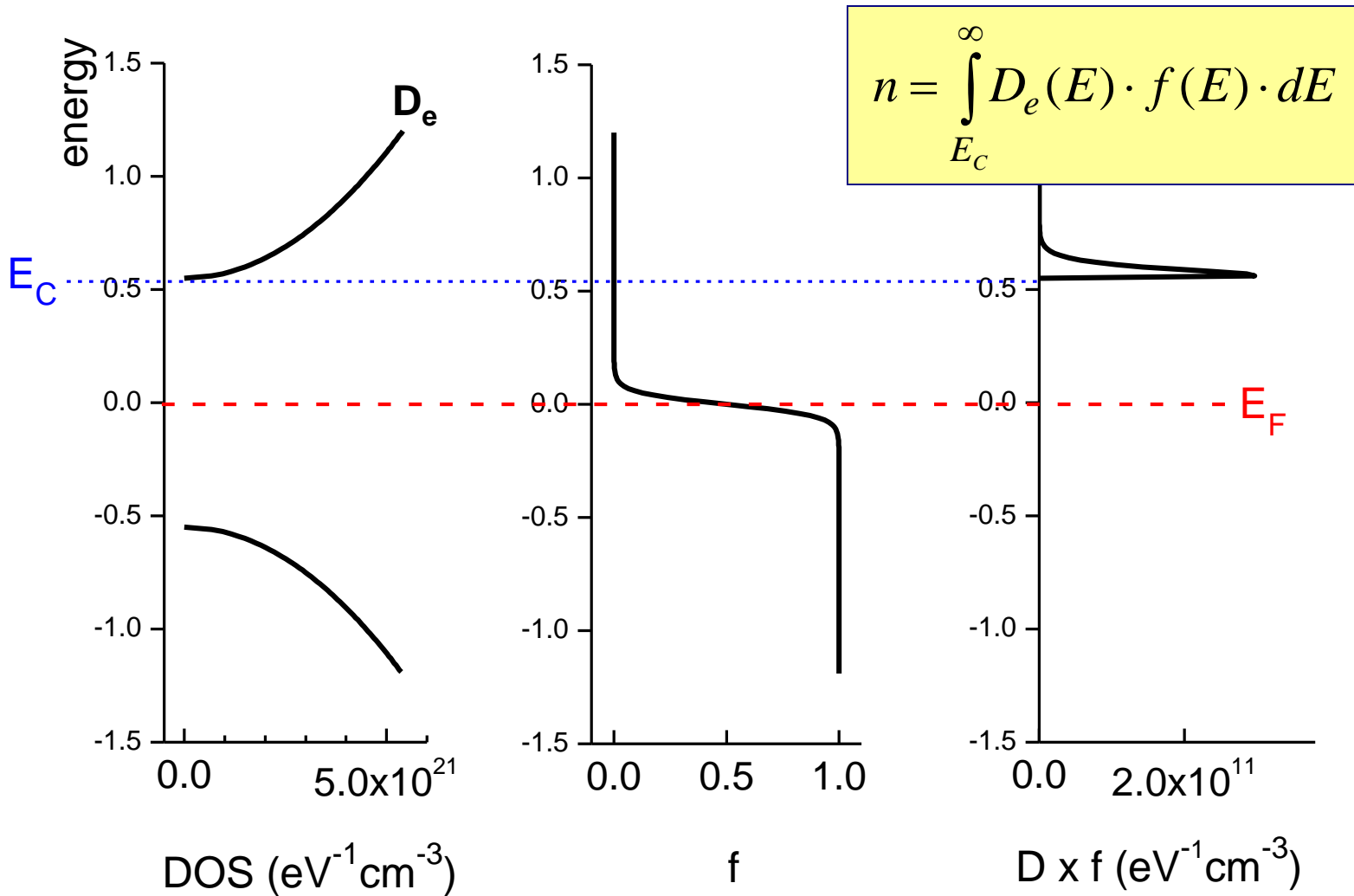
$E_F$  – Fermi-energy

distribution function following  
from the nature of the electron

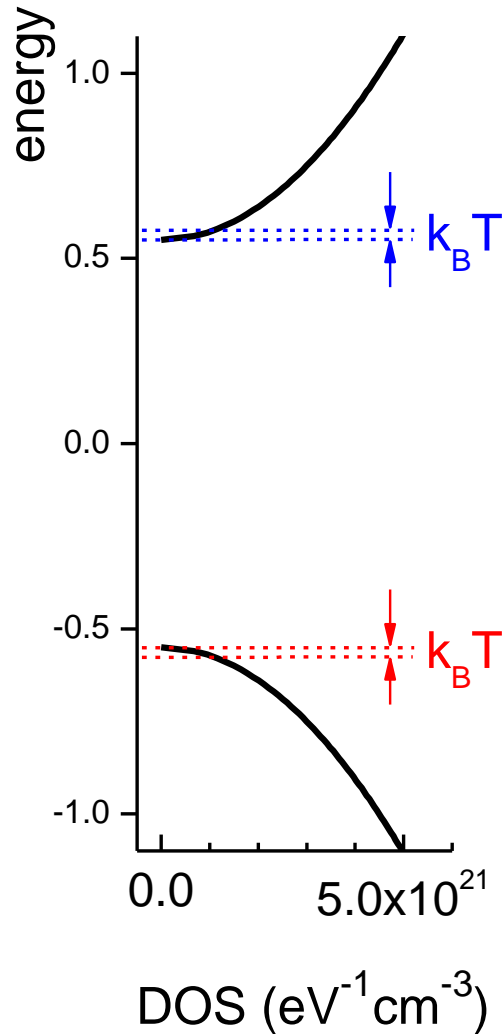




# The product of $D_e(E)$ and $f(E)$



# The effective DOS



$$N_C \gg \int_{E_C}^{E_C + k_B T} D_e(E) \times dE$$

$$N_V \gg \int_{E_V - k_B T}^{E_V} D_h(E) \times dE$$

$$N_{C(V)} = 2 \cdot \left( 2\pi \cdot k_B \cdot T \cdot \frac{m_{e(h)}^*}{h^2} \right)^{3/2}$$

	$N_C (cm^{-3})$	$N_V (cm^{-3})$
Si	$2.4 \cdot 10^{19}$	$1 \cdot 10^{19}$
Ge	$1 \cdot 10^{19}$	$6 \cdot 10^{18}$
GaAs	$5 \cdot 10^{17}$	$7 \cdot 10^{18}$

# Densities of free charge carriers

$$n = N_C \cdot \exp\left(-\frac{E_C - E_{Fn}}{k_B \cdot T}\right)$$

$$p = N_V \cdot \exp\left(-\frac{E_{Fp} - E_V}{k_B \cdot T}\right)$$

Boltzmann statistics

# Effective density of states

$$N_C = 2 \cdot \left( 2\pi \cdot k_B \cdot T \cdot \frac{m_e^*}{h^2} \right)^{3/2}$$

$$N_C = 2.5 \left( \frac{m_e^*}{m_e} \right)^{3/2} \left( \frac{T}{300K} \right)^{3/2} \times 10^{19} \text{ cm}^{-3}$$

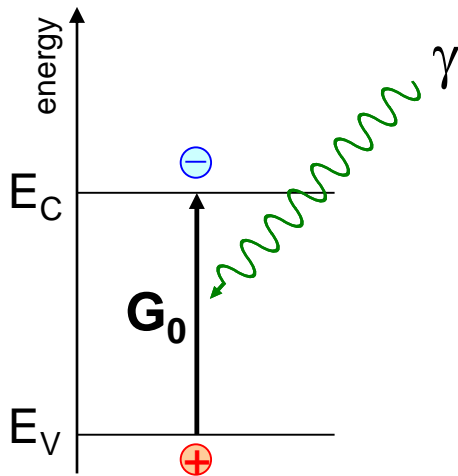
$$N_V = 2 \cdot \left( 2\pi \cdot k_B \cdot T \cdot \frac{m_h^*}{h^2} \right)^{3/2}$$

$$N_V = 2.5 \left( \frac{m_h^*}{m_e} \right)^{3/2} \left( \frac{T}{300K} \right)^{3/2} \times 10^{19} \text{ cm}^{-3}$$

1. Free electron and hole densities
2. Intrinsic carrier concentration
3. Doping of semiconductors
4. Equilibrium and excess charge carriers

# Re-absorption of black body irradiation

semiconductor emits black body irradiation  
semiconductor absorbs its own black body irradiation for  $h\nu \geq E_g$



$$\gamma \rightarrow e^- + h^+$$

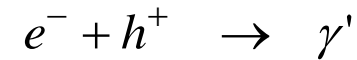
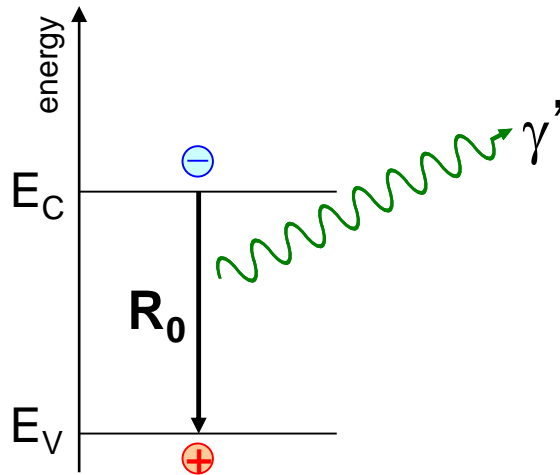
$G_0$  thermal generation rate  
of free charge carriers

# Recombination of free charge carriers

free carrier concentrations can not increase infinitely



limitation by recombination



$R_0$  recombination rate

$$R_0 = B \cdot n_0 \cdot p_0$$

$R_0$  is proportional to  $n_0$  and  $p_0$   
 $B$  recombination rate constant

# Thermal equilibrium

$$G_0 = R_0$$



$$n_0 \cdot p_0 = \frac{G_0}{B} = \text{const}$$

$$n_0 \cdot p_0 = N_C \cdot N_V \cdot \exp\left(-\frac{E_C - E_{Fn} + E_{Fp} - E_V}{k_B \cdot T}\right)$$

$$n_0 \cdot p_0 = N_C \cdot N_V \cdot \exp\left(-\frac{E_g}{k_B \cdot T}\right) \cdot \exp\left(-\frac{E_{Fp} - E_{Fn}}{k_B \cdot T}\right)$$

$E_{Fn}$  and  $E_{Fp}$  can change with the free carrier concentration



$$E_{Fn} = E_{Fp}$$



# Intrinsic carrier concentration

in thermal equilibrium  $n_0 \cdot p_0 = \text{const} \equiv n_i^2$

$n_i$  intrinsic carrier concentration

$$n_i = \sqrt{n_0 \cdot p_0}$$

# Intrinsic carrier concentration

in thermal equilibrium

$$n_i = \sqrt{n_0 \cdot p_0}$$

from Boltzmann statistics

$$n_0 = N_C \cdot \exp\left(-\frac{E_C - E_F}{k_B \cdot T}\right)$$

$$p_0 = N_V \cdot \exp\left(-\frac{E_F - E_V}{k_B \cdot T}\right)$$

# Intrinsic carrier concentration

in thermal equilibrium

$$n_i = \sqrt{n_0 \cdot p_0}$$

Boltzmann statistics

$$n_0 = N_C \cdot \exp\left(-\frac{E_C - E_F}{k_B \cdot T}\right)$$

$$p_0 = N_V \cdot \exp\left(-\frac{E_F - E_V}{k_B \cdot T}\right)$$

$$n_i = \sqrt{N_V \cdot N_C} \cdot \exp\left(-\frac{E_g}{2 \cdot k_B \cdot T}\right)$$

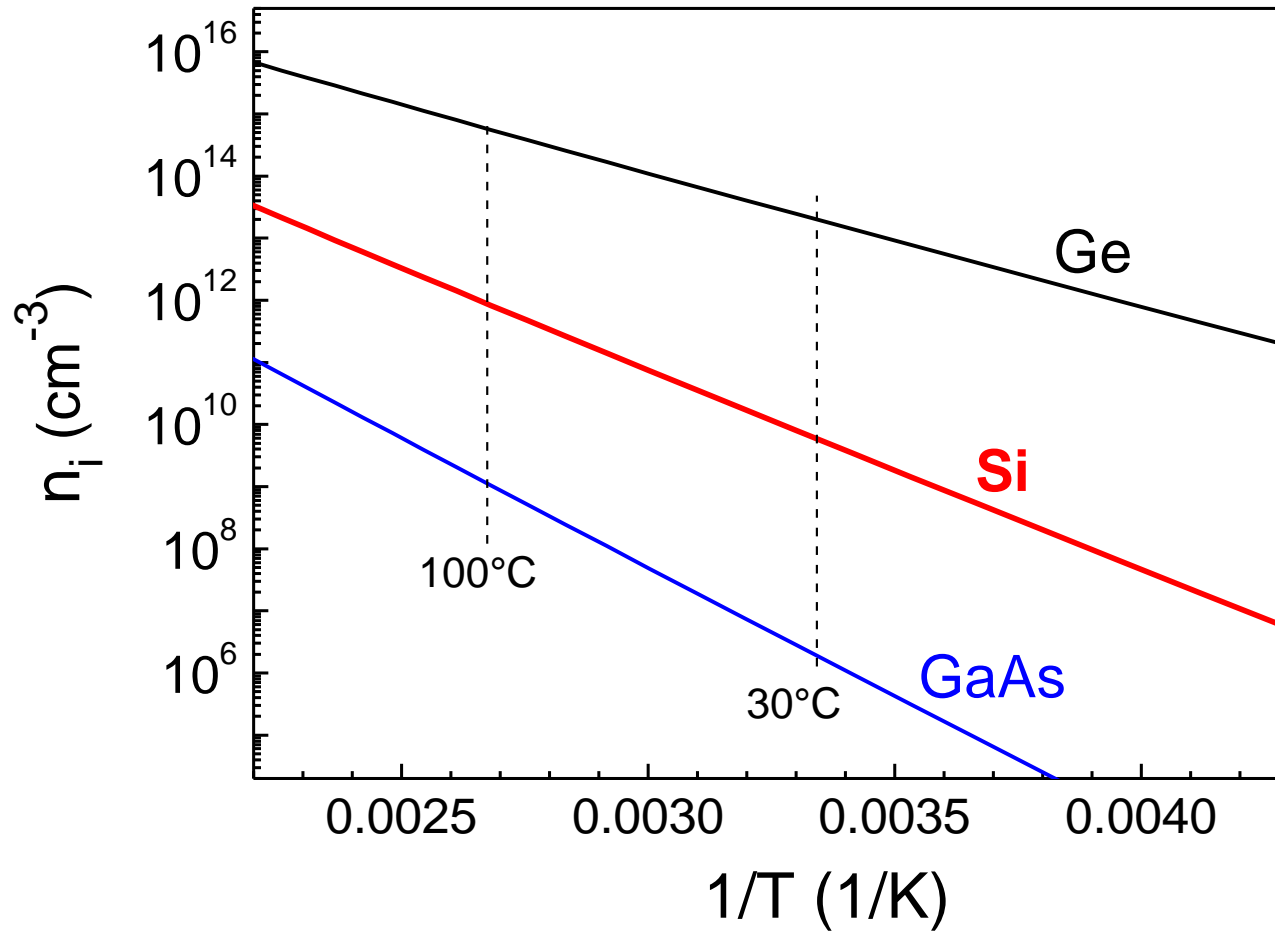
intrinsic semiconductor

$$n_i = n_0 = p_0$$

intrinsic Fermi-level

$$E_F = E_i = \frac{E_C + E_V}{2} + \frac{3}{4} \cdot k_B \cdot T \cdot \ln \frac{m_h^*}{m_e^*}$$

# Temperature dependence of the intrinsic carrier concentration



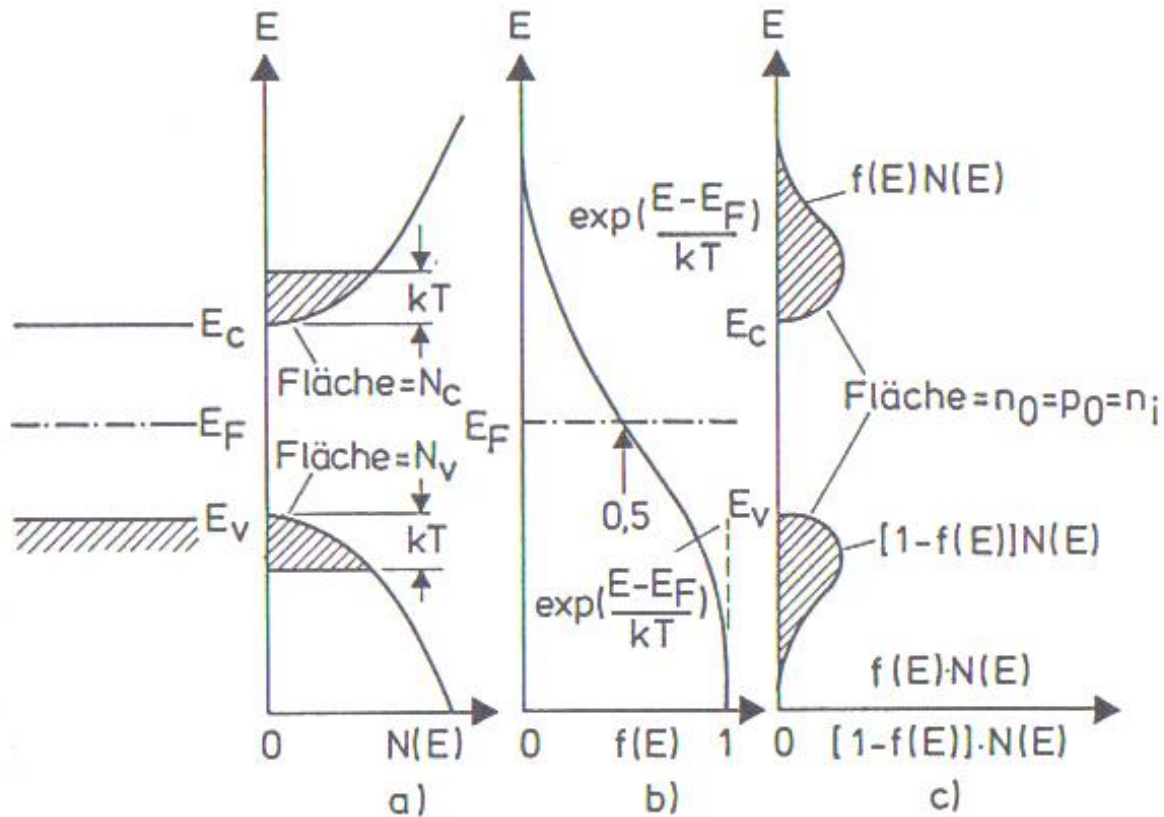
# Intrinsic carrier concentrations at room temperature

$$n_i(\text{Si}) \approx 10^{10} \text{ cm}^{-3}$$

$$n_i(\text{GaAs}) \approx 10^6 \text{ cm}^{-3}$$

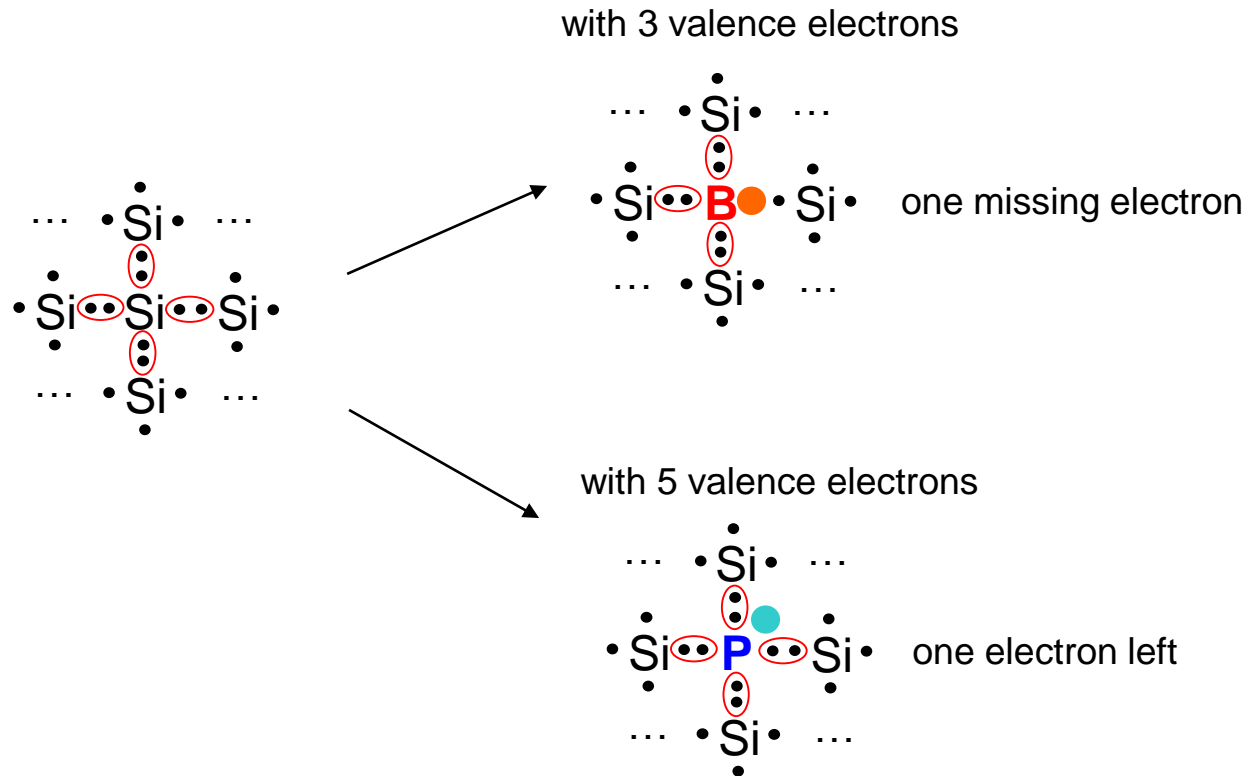
$$n_i(\text{Ge}) \approx 10^{13} \text{ cm}^{-3}$$

# Intrinsic semiconductor



1. Free electron and hole densities
2. Intrinsic carrier concentration
3. Doping of semiconductors
4. Equilibrium and excess charge carriers

# Impurity in a crystal lattice



There are impurities which can be incorporated into the lattice replacing a host atom. importance of valence



# Doping of semiconductors

Incorporation of impurity atoms into the crystal lattice of the semiconductor.

Impurity atoms have different valence in comparison to the atoms of the host lattice.

## n-type doping

The impurity atom has more valence electrons than needed for bonding.

Impurity atom giving an electron is named **donor**.

Un-bonded electron can be excited into the conduction band.

creation of a mobile electron in the conduction band

## p-type doping

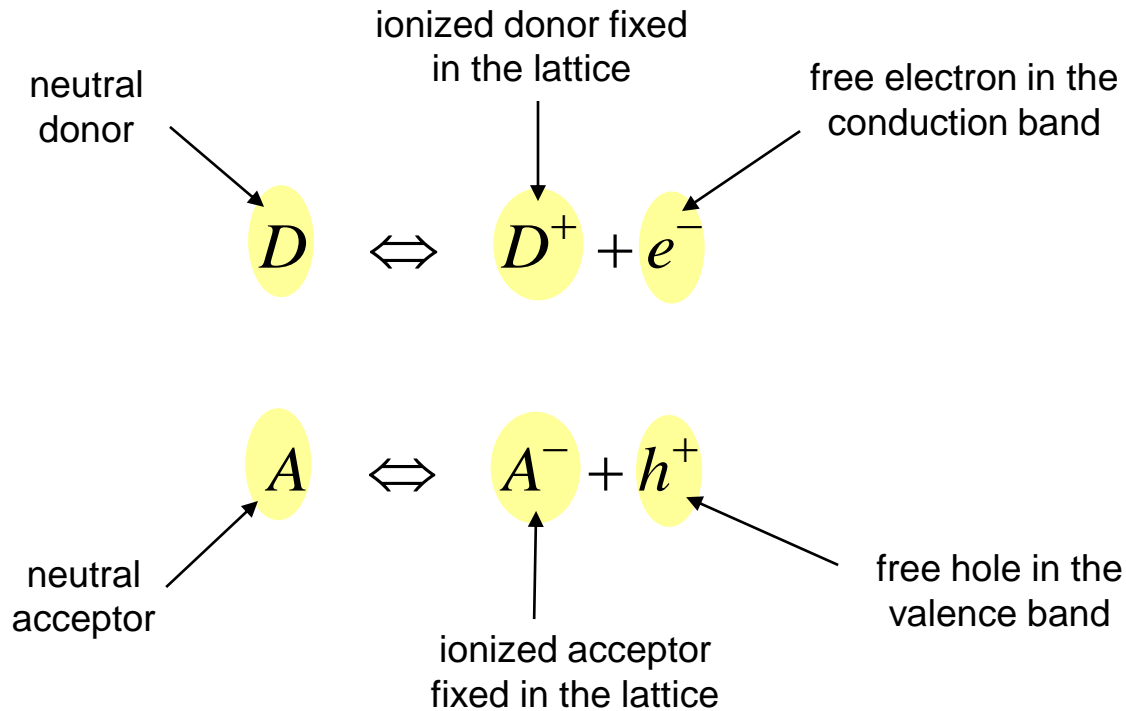
The impurity atom has less valence electrons than needed for bonding.

Impurity atom taking an electron is named **acceptor**.

Electron needed for bonding can be taken from the valence band.

creation of a mobile hole in the valence band

# Free charge carriers and fixed ionized impurity



charge neutrality:

n-type

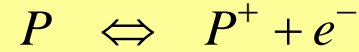
$$N_D^+ = n_0$$

p-type

$$N_A^- = p_0$$

# Doping of Si

donors: P, As, ...  
acceptors: B, Al, ...



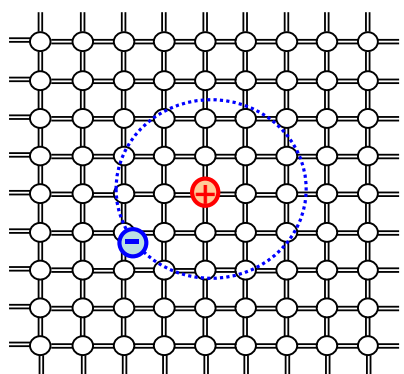
## PERIODIC TABLE OF ELEMENTS

AVAILABLE AS NANOCRYSTALLINE METAL

AVAILABLE AS NANOCRYSTALLINE METAL OXIDES

<b>IA</b>												<b>VIIIB</b>					
<b>H 1</b> HYDROGEN 0.0899 1.0794											<b>He 2</b> HELIUM 0.1785 4.003						
<b>IIA</b>												<b>IIIA</b>	<b>IVA</b>	<b>VA</b>	<b>VIA</b>	<b>VIIA</b>	<b>Ne 10</b>
<b>Li 3</b> LITHIUM 0.534 6.941	<b>Be 4</b> BERYLLIUM 1.848 9.012184											<b>B 5</b> BORON 2.34 10.811	<b>C 6</b> CARBON 2.25 12.01	<b>N 7</b> NITROGEN 1.25 14.006	<b>O 8</b> OXYGEN 1.429 15.999	<b>F 9</b> FLUORINE 1.696 18.998	<b>Ne 10</b> NEON 0.900 20.179
<b>Na 11</b> SODIUM 0.971 22.989768		<b>Mg 12</b> MAGNESIUM 1.738 24.305		<b>IIIB</b>	<b>IVB</b>	<b>VB</b>	<b>VIB</b>	<b>VIIB</b>	<b>VIII</b>	<b>IB</b>	<b>IIB</b>	<b>Ga 31</b> GALLIUM 2.6889 69.723	<b>Ge 32</b> GERMANIUM 2.33 72.61	<b>As 33</b> ARSENIC 1.82 74.921	<b>Se 34</b> SELENIUM 2.07 78.96	<b>Br 35</b> BROMINE 3.214 79.904	<b>Kr 36</b> KRYPTON 1.783 83.94
<b>K 19</b> POTASSIUM 0.862 39.0983	<b>Ca 20</b> CALCIUM 1.55 40.078	<b>Sc 21</b> SCANDIUM 2.989 44.956	<b>Ti 22</b> TITANIUM 4.54 47.88	<b>V 23</b> VANADIUM 6.11 50.941	<b>Cr 24</b> CHROMIUM 7.19 51.996	<b>Mn 25</b> MANGANESE 7.44 54.938	<b>Fe 26</b> IRON 7.874 55.847	<b>Co 27</b> COBALT 8.90 58.933	<b>Ni 28</b> NICKEL 8.902 58.69	<b>Cu 29</b> COPPER 8.90 63.546	<b>Zn 30</b> ZINC 7.133 65.39	<b>Ga 31</b> GALLIUM 5.904 69.723	<b>Ge 32</b> GERMANIUM 5.323 72.61	<b>As 33</b> ARSENIC 5.727 74.921	<b>Se 34</b> SELENIUM 4.79 78.96	<b>Br 35</b> BROMINE 3.12 79.904	<b>Kr 36</b> KRYPTON 3.733 83.80
<b>Rb 37</b> RUBIDIUM 1.532 85.467	<b>Sr 38</b> STRONTIUM 2.54 87.62	<b>Y 39</b> YTTORIUM 4.469 88.905	<b>Zr 40</b> ZIRCONIUM 6.506 91.224	<b>Nb 41</b> NIObIUM 8.57 92.906	<b>Mo 42</b> MOLYBDENUM 10.22 95.94	<b>Tc 43</b> TECHNETIUM 11.50 (99)	<b>Ru 44</b> RUTHENIUM 12.46 101.07	<b>Rh 45</b> RHODIUM 12.41 102.905	<b>Pd 46</b> PALLADIUM 12.02 106.42	<b>Ag 47</b> SILVER 10.50 107.868	<b>Cd 48</b> CADMIUM 8.65 112.411	<b>In 49</b> INDIUM 7.31 114.82	<b>Sn 50</b> TIN 7.31 118.71	<b>Sb 51</b> ANTIMONY 6.691 121.75	<b>Te 52</b> TELLURIUM 6.24 127.60	<b>I 53</b> IODINE 4.93 126.904	<b>Xe 54</b> XENON 5.89 131.29
<b>Cs 55</b> CESIUM 1.873 132.905	<b>Ba 56</b> BARIUM 3.5 137.327	<b>La 57</b> LANTHANIUM 6.145 138.905	<b>Hf 72</b> HAFNIUM 13.31 178.49	<b>Ta 73</b> TANTALUM 16.654 180.947	<b>W 74</b> TUNGSTEN 19.3 183.85	<b>Re 75</b> RHENIUM 21.02 186.207	<b>Os 76</b> OSMIUM 22.57 190.2	<b>Ir 77</b> IRIDIUM 22.42 192.22	<b>Pt 78</b> PLATINUM 21.45 195.08	<b>Au 79</b> GOLD 19.31 196.96	<b>Hg 80</b> MERCURY 13.46 200.59	<b>Tl 81</b> THALLIUM 11.85 204.383	<b>Pb 82</b> LEAD 11.35 207.2	<b>Bi 83</b> BISMUTH 9.747 208.98	<b>Po 84</b> POLONIUM 9.32 (209)	<b>At 85</b> ASTATINE (210)	<b>Rn 86</b> RADON 9.73 (222)
<b>Fr 87</b> FRANCIUM (223)	<b>Ra 88</b> RADIUM 5 226.025	<b>Ac 89</b> ACTINIUM 10.07 227.027	<b>Rf 104</b> (261)	<b>Ha 105</b> (262)	<b>Sg 106</b> (263)	<b>Ns 107</b> (262)	<b>Hs 108</b> (262)	<b>Mt 109</b> (262)									
<b>*LANTHANIDE SERIES</b>			<b>Ce 58</b> CERIUM 6.77 140.115	<b>Pr 59</b> PRAESESODIUM 6.773 140.908	<b>Nd 60</b> NEODYMIUM 7.008 144.24	<b>Pm 61</b> PROMETHIUM 7.22 (145)	<b>Sm 62</b> SAMARIUM 7.52 150.36	<b>Eu 63</b> EUROPIUM 5.24 151.965	<b>Gd 64</b> GADOLINIUM 7.901 157.25	<b>Tb 65</b> TERBIUM 8.230 158.925	<b>Dy 66</b> DYSPROSIUM 8.551 162.50	<b>Ho 67</b> HOLMIUM 8.795 164.930	<b>Er 68</b> ERBIUM 9.066 167.26	<b>Tm 69</b> THULIUM 9.321 168.934	<b>Yb 70</b> YTTERIUM 6.966 173.04	<b>Lu 71</b> LUTETIUM 9.841 174.967	
<b>**ACTINIDE SERIES</b>			<b>Th 90</b> THORIUM 11.72 232.038	<b>Pa 91</b> PROTOACTINIUM 15.37 231.036	<b>U 92</b> URANIUM 18.95 238.029	<b>Np 93</b> NEPTUNIUM 20.25 237.04	<b>Pu 94</b> PLUTONIUM 19.84 (244)	<b>Am 95</b> AMERICIUM 13.67 (243)	<b>Cm 96</b> CURIUM 13.51 (247)	<b>Bk 97</b> BERKELIUM 14.00	<b>Cf 98</b> CALIFORNIUM	<b>Es 99</b> EINSTEINIUM	<b>Fm 100</b> FERMIUM	<b>Md 101</b> MENDELEVIUM (259)	<b>No 102</b> NOBELIUM (260)	<b>Lr 103</b> LAWRENCIUM (260)	

# Hydrogen model of doping



ionization energy

$$E_{ion} \propto \frac{m^*}{\epsilon^2}$$

Bohr-radius

$$a_0 \propto \frac{\epsilon}{m^*}$$

$$E_{ion}(H - atom) = 13.56 \text{ eV}$$

$$a_0(H - atom) = 0.051 \text{ nm}$$

$$\epsilon_{Si} = 11.9$$

estimation of the ionization or binding energy

$$E_{ion} = E_{ion}(H - atom) \cdot \frac{m_{e(Si, GaAs, \dots)}^*}{m_e} \cdot \frac{\epsilon_{vac}^2}{\epsilon_{Si, GaAs, \dots}^2}$$

$$E_{ion} = 6 \text{ meV (GaAs)} \dots \approx 50 \text{ meV (Si)}$$



all impurities are ionized at room temperature

$$N_D \approx N_D^+ = n_0$$

$$N_A \approx N_A^- = p_0$$

# Energy levels of some impurities in Si

donors

	$E_C - E_D$ (eV)
Sb	0.039
P	0.044
As	0.049
Bi	0.069

acceptors

	$E_V + E_A$ (eV)
B	0.046
Al	0.057
Ga	0.065
In	0.16

# Minority and majority charge carriers

majority charge carriers:

electrons in n-type  
hole in p-type

minority carriers:

holes in n-type  
electrons in p-type

$$n_i^2 = n_0 \cdot p_0$$

Example: p-Si,  $N_A = 10^{16} \text{ cm}^{-3}$

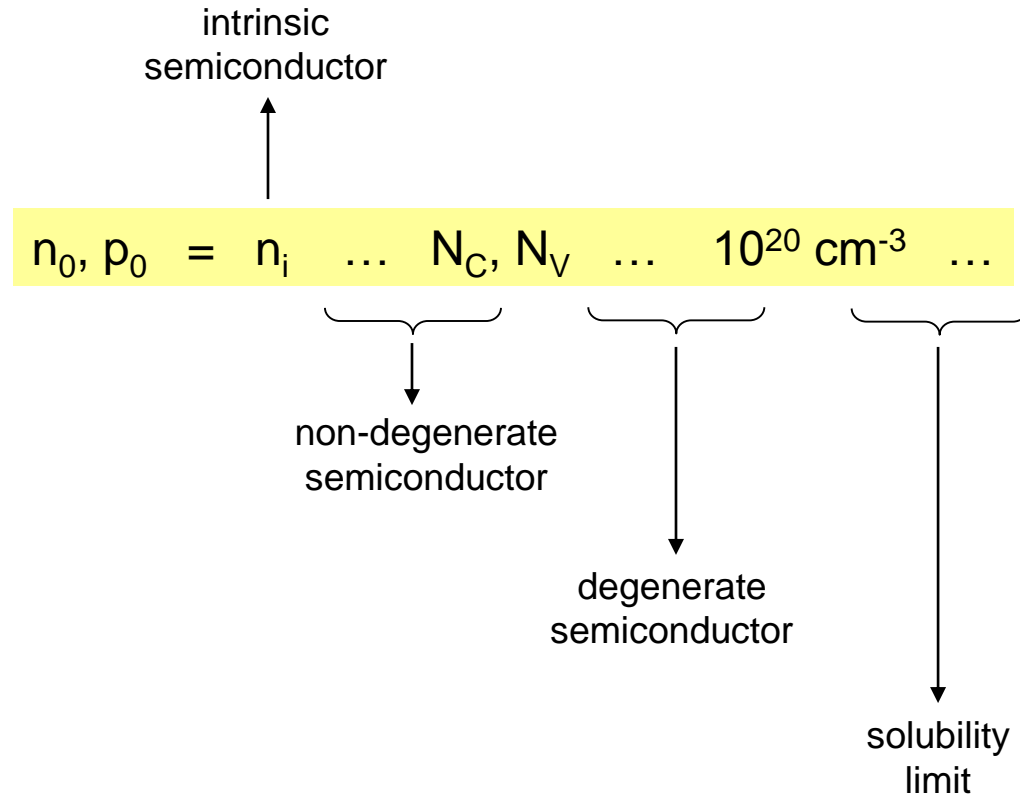
$n_i(\text{Si}) \approx 10^{10} \text{ cm}^{-3}$

majority:  $p_0 = 10^{16} \text{ cm}^{-3}$

minority:  $n_0 = 10^4 \text{ cm}^{-3}$

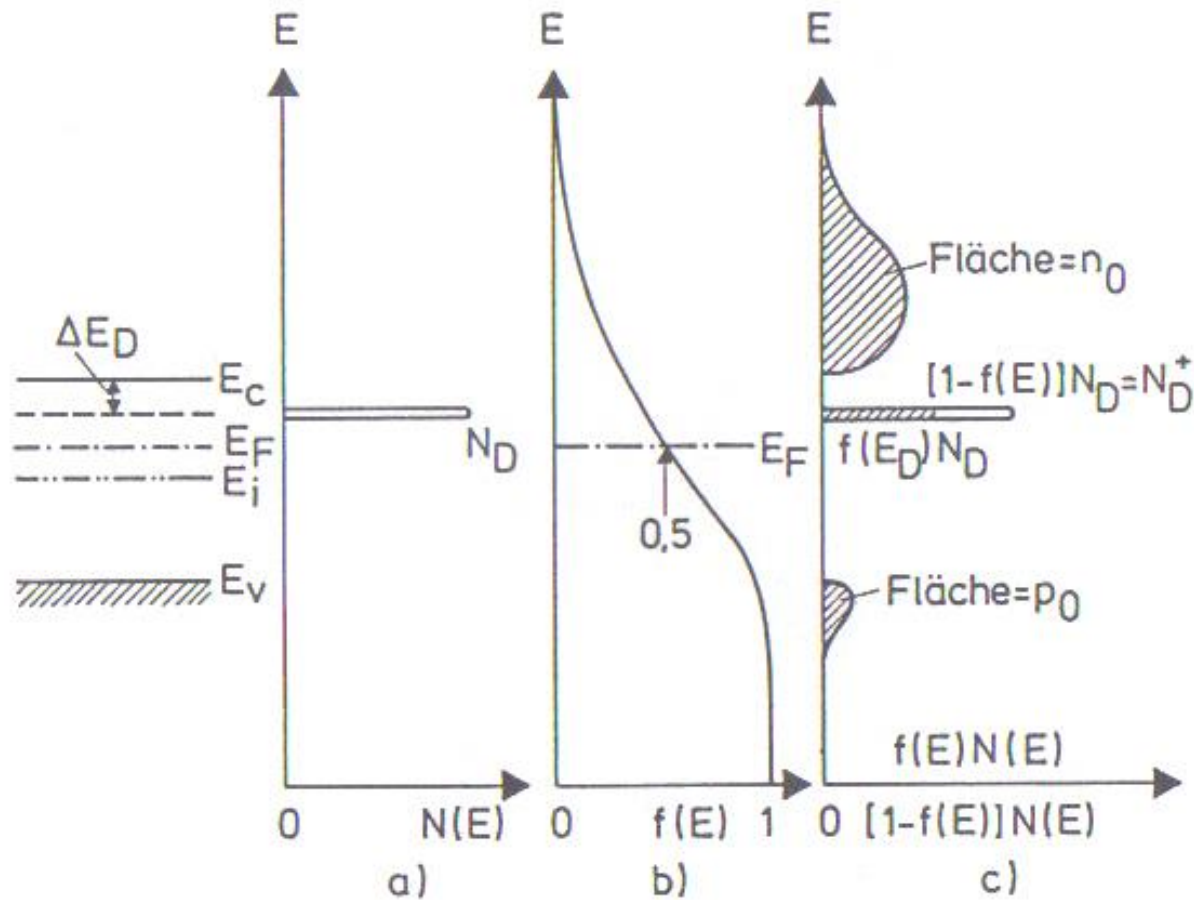
„law of mass action“  
(it is sufficient to know one  
density of charge carriers)

# Doping range of semiconductors



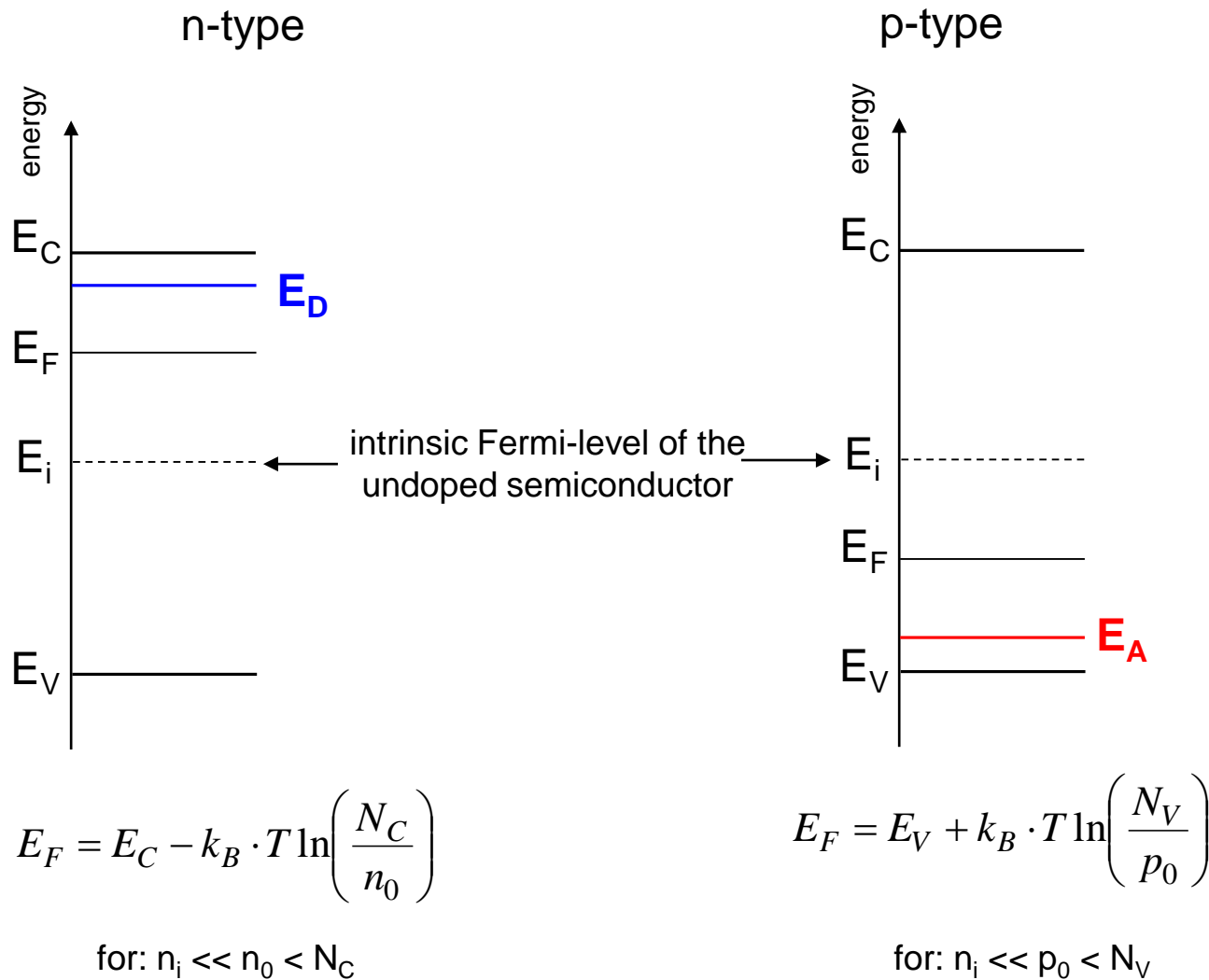
for Si from minority to majority:  $\approx 20$  orders of magnitude  
huge change of conductivity (about 10 orders of magnitude for majority carriers)

# Doped semiconductor





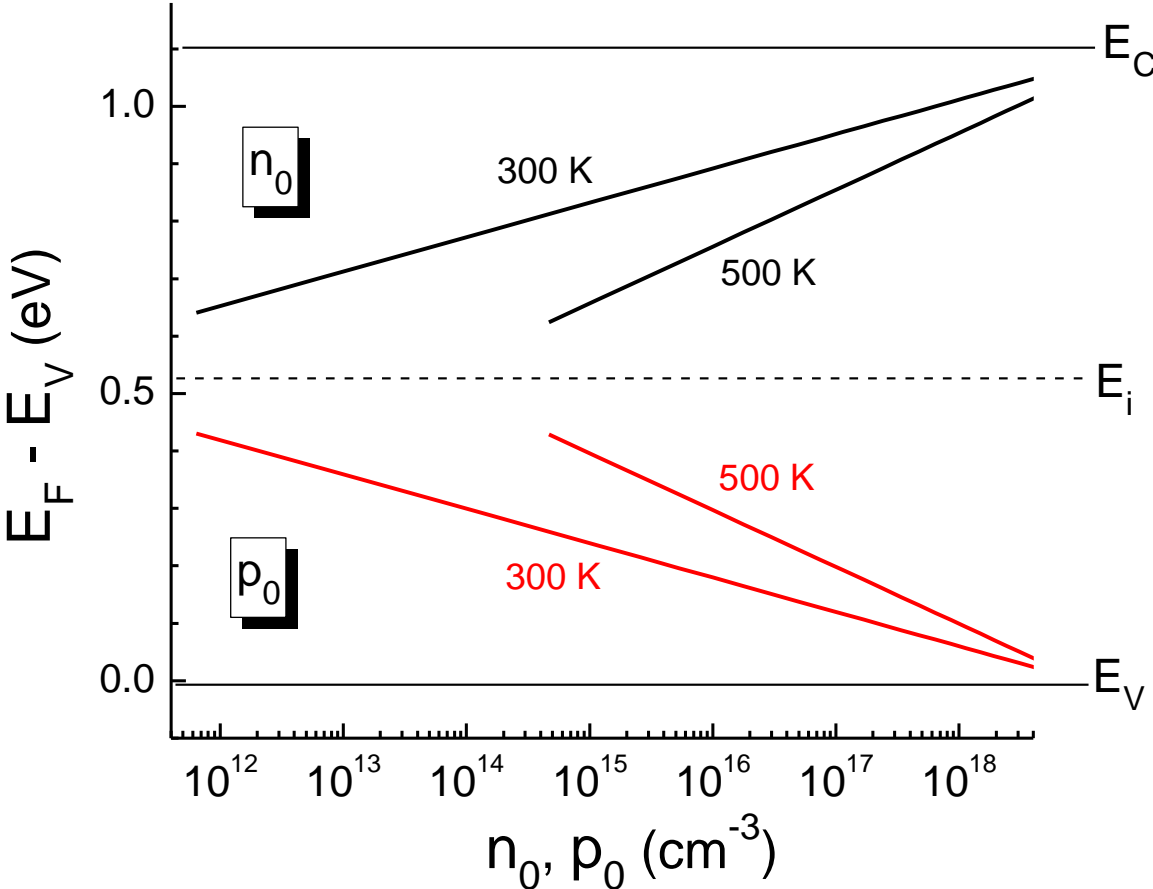
# Energetic levels in doped semiconductors



# Formulaes for the doped semiconductor

	$n_e$	$n_h$	$\varepsilon_F$
n-Leiter	$n_e \approx n_D$	$n_h = \frac{n_i^2}{n_e} = \frac{n_i^2}{n_D}$	$\varepsilon_C - kT \ln \frac{N_C}{n_D}$
p-Leiter	$n_e = \frac{n_i^2}{n_h} = \frac{n_i^2}{n_A}$	$n_h \approx n_A$	$\varepsilon_V + kT \ln \frac{N_V}{n_A}$

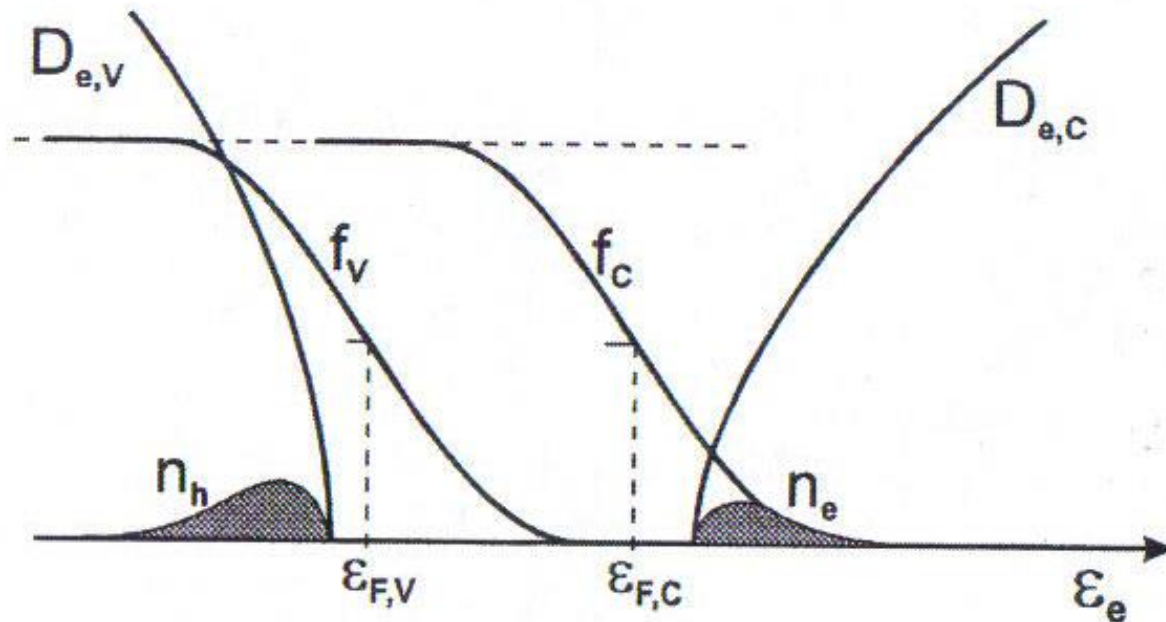
# Doping dependence of the Fermi-level



(temperature dependence of  $E_g$  has not been considered)

1. Free electron and hole densities
2. Intrinsic carrier concentration
3. Doping of semiconductors
4. Equilibrium and excess charge carriers

# Quasi-Fermi-levels under illumination



**Abb. 3.12** Im belichteten Halbleiter wird die Besetzung von Leitungsband  $C$  und Valenzband  $V$  von verschiedenen Fermi-Verteilungen  $f_c$  und  $f_v$  geregelt.

thermal equilibrium

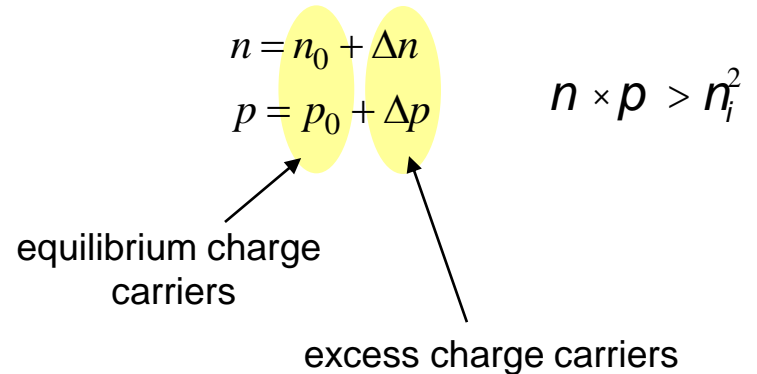
$$E_F = E_{F0} = E_{Fn} = E_{Fp}$$

$$n_i^2 = n_0 \cdot p_0$$

$$\begin{aligned} E_{F0} &= E_C - k_B \cdot T \cdot \ln\left(\frac{N_C}{n_0}\right) \\ &= E_V + k_B \cdot T \cdot \ln\left(\frac{N_V}{p_0}\right) \end{aligned}$$

under illumination  
(or non-equilibrium)

$$E_{Fn} \neq E_{Fp}$$



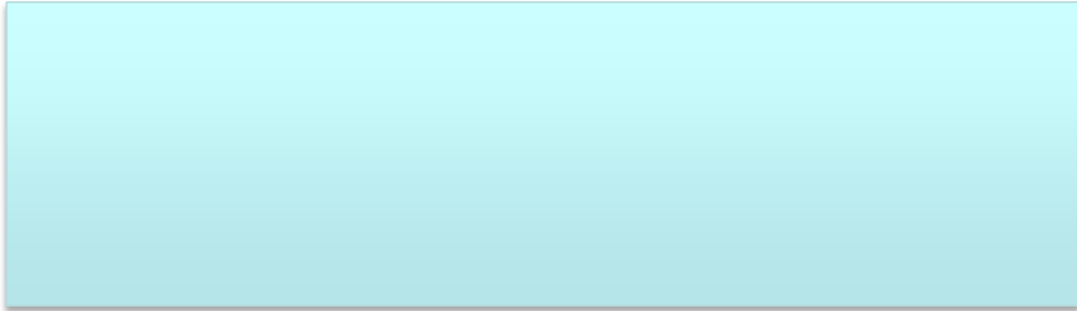
$$\begin{aligned} E_{Fn} &= E_C - k_B \cdot T \cdot \ln\left(\frac{N_C}{n_0 + \Delta n}\right) \\ E_{Fp} &= E_V + k_B \cdot T \cdot \ln\left(\frac{N_V}{p_0 + \Delta p}\right) \end{aligned}$$

separation into quasi Fermi-levels in non-equilibrium

# Ideal charge-selective contact

$$q \cdot V_{OC} = E_{Fn} - E_{Fp}$$

no losses of potential energy at contacts


$$= k_B \cdot T \cdot \ln \left( \frac{n_0 + \Delta n}{n_0} \cdot \frac{p_0 + \Delta p}{p_0} \right)$$

# Origin of $V_{OC}$

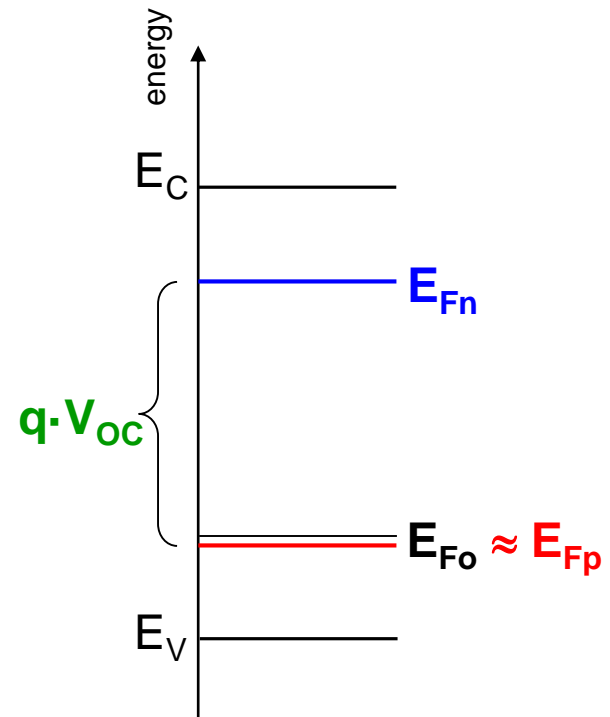
$$V_{OC} = \frac{k_B \cdot T}{q} \cdot \ln \left( \frac{n_0 + \Delta n}{n_0} \cdot \frac{p_0 + \Delta p}{p_0} \right)$$

$$V_{OC} = \frac{k_B \cdot T}{q} \cdot \ln \left( \frac{\Delta n}{n_0} \right)$$

example: p-type Si

$$n_0 + \Delta n \approx \Delta n$$

$$p_0 + \Delta p \approx p_0$$



Concentration ratio of photo generated and minority charge carriers determines  $V_{OC}$ .



## Example: p-Si under illumination

typical doping range of the base:  $p_0 = 10^{16} \text{ cm}^{-3}$   
 $n_0 = 10^4 \text{ cm}^{-3}$

typical concentration of photo generated charge carriers  
 $\Delta n = \Delta p = 10^{15} \text{ cm}^{-3}$

$$E_{F0} = E_V + 0.179 \text{ eV}$$

$$E_{Fn} = E_C - 0.26 \text{ eV} = E_g + E_V - 0.26 \text{ eV} = E_V + 0.857 \text{ eV}$$

$$E_{Fp} = E_V + 0.177 \text{ eV}$$

$$E_{F0} - E_{Fp} = 0.002 \text{ eV}$$

$$E_{Fn} - E_{F0} = 0.680 \text{ eV}$$

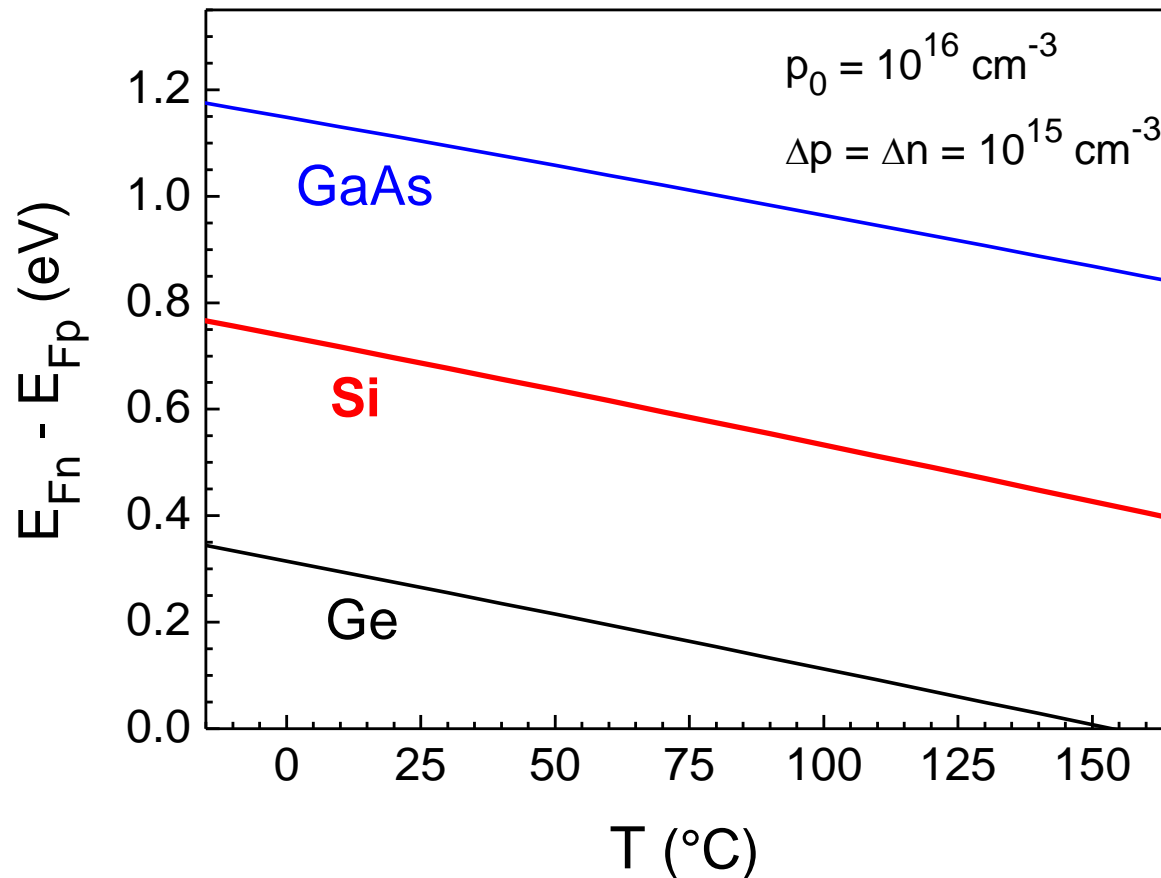
$$E_{Fn} - E_{Fp} = 0.678 \text{ eV}$$

$$E_{F0} - E_{Fp} \ll E_{Fn} - E_{F0}$$

# Temperature dependent separation of quasi Fermi-levels

$$n_i^2 = N_V \cdot N_C \cdot \exp\left(-\frac{E_g}{kT}\right)$$

$$E_{Fn} - E_{Fp} = kT \cdot \ln\left(\frac{(n_0 + \Delta n) \cdot (p_0 + \Delta p)}{n_i^2}\right)$$



Strong temperature dependence of  $n_i$  is the reason for the decrease of  $V_{OC}$  with increasing temperature.