

2a) classical torque: $\vec{T} = \vec{\mu} \wedge \vec{B}$

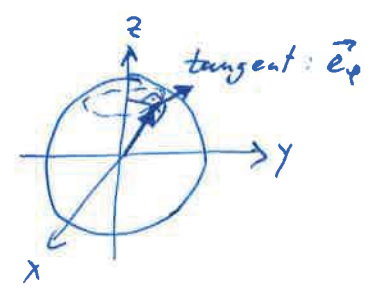
explicitly: in cartesian coordinates we have $\vec{\mu} = \mu \begin{pmatrix} \sin(\theta) \cdot \cos(\varphi) \\ \sin(\theta) \cdot \sin(\varphi) \\ \cos(\theta) \end{pmatrix}$

$$\Rightarrow \vec{T} = \vec{\mu} \wedge \vec{B} = \mu \begin{pmatrix} \sin(\theta) \cos(\varphi) \\ \sin(\theta) \sin(\varphi) \\ \cos(\theta) \end{pmatrix} \wedge \begin{pmatrix} 0 \\ 0 \\ B \end{pmatrix} = \mu \cdot B \cdot \begin{pmatrix} \sin(\theta) \sin(\varphi) \\ -\sin(\theta) \cos(\varphi) \\ 0 \end{pmatrix}$$

$$= \mu \cdot B \cdot \sin(\theta) \cdot \begin{pmatrix} \sin(\varphi) \\ -\cos(\varphi) \\ 0 \end{pmatrix}$$

$= -\vec{e}_\varphi$: unit vector for changes in φ in spherical coordinates

$\Rightarrow \underline{\underline{\vec{T} = \mu B \sin(\theta) \cdot \vec{e}_\varphi}}$

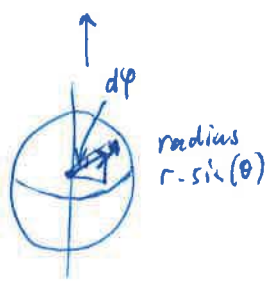
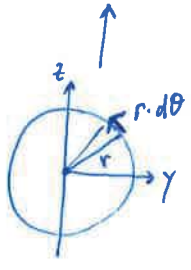


b) goal: i) derive differential equation for $\vec{L}(t)$, i.e. obtain $\frac{\partial \vec{L}}{\partial t}$
 ii) solve differential equation

i) In polar coordinates the change in a vector \vec{s} when the coordinates are changed by dr (radius), $d\theta$ and $d\varphi$ (angles) is

$$d\vec{s} = \vec{e}_r \cdot dr + \vec{e}_\theta \cdot r \cdot d\theta + \vec{e}_\varphi \cdot r \cdot \sin(\theta) \cdot d\varphi$$

↑
radial unit vector



$$\Rightarrow \frac{d\vec{s}}{dt} = \vec{e}_r \cdot \frac{dr}{dt} + \vec{e}_\theta \cdot r \cdot \frac{d\theta}{dt} + \vec{e}_\varphi \cdot r \cdot \sin(\theta) \cdot \frac{d\varphi}{dt}$$

or for \vec{L} : $\underline{\underline{\frac{d\vec{L}}{dt} = \vec{e}_r \cdot \frac{dL}{dt} + \vec{e}_\theta \cdot L \cdot \frac{d\theta}{dt} + \vec{e}_\varphi \cdot L \cdot \sin(\theta) \cdot \frac{d\varphi}{dt}}}$

$\stackrel{!}{=} \vec{T} = -\mu B \sin(\theta) \cdot \vec{e}_\varphi$ (Differential equation for \vec{L})
 ↑
by definition

with $\mu = -\frac{2m}{\hbar} \cdot \mu$

ii) solve differential equation(s):
 • radial component: $\frac{dL}{dt} \stackrel{!}{=} 0 \Rightarrow \underline{\underline{L = const.}}$
 • \vec{e}_θ -component: $L \cdot \frac{d\theta}{dt} \stackrel{!}{=} 0, L \neq 0 \Rightarrow \underline{\underline{\theta = const.}}$

• \vec{e}_φ -component: $L \cdot \sin(\theta) \frac{d\varphi}{dt} \stackrel{!}{=} \mu B \cdot \sin(\theta)$

$$\frac{d\varphi}{dt} = -\frac{\mu B}{L} = -\frac{\mu B}{-\frac{2m}{g_L e} \mu} = \frac{g_L e}{2m} \cdot B$$

$$\Rightarrow \underline{\underline{\varphi(t) = \frac{g_L e}{2m} \cdot B \cdot t}}$$

(3)

c) Since L, θ are constant and $\varphi \propto t$ we obtain

$$\vec{L}(t) = L \cdot \begin{pmatrix} \sin(\theta) \cos(\varphi) \\ \sin(\theta) \sin(\varphi) \\ \cos(\theta) \end{pmatrix}, \text{ i.e. } L_z = \text{const.}, \quad L_x = L \cdot \sin(\theta) \cdot \cos(\varphi) = L_{\parallel} \cdot \cos(\varphi(t))$$

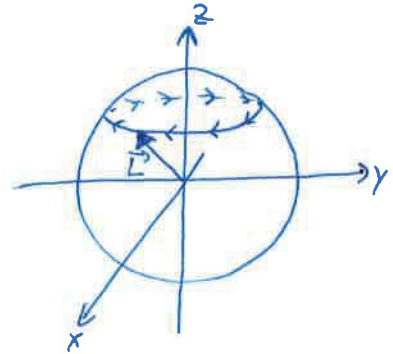
$$L_y = L \cdot \sin(\theta) \cdot \sin(\varphi) = L_{\parallel} \cdot \sin(\varphi(t))$$

which describes a circle with radius $L_{\parallel} = L \cdot \sin(\theta)$.

Rotation frequency: full period $\hat{=} \varphi(\Delta t) = 2\pi = \frac{g_L e}{2m} B \cdot \Delta t$

$$\Rightarrow \underline{\underline{\omega = \frac{2\pi}{\Delta t} = \frac{g_L e}{2m} \cdot B}}$$

(Larmor frequency)



Larmor - precession :

$$\psi_+^z = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \psi_-^z = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \text{eigenstates of } S_z$$

$$H = -\gamma \vec{B} \cdot \vec{S} \quad E_+ = -\frac{\hbar}{2} \gamma B \quad E_- = \frac{\hbar}{2} \gamma B$$

time evolution

usually: $\psi(t) = \sum_n c_n e^{-iE_n t/\hbar} \psi_n$

$$\psi(t) = a \psi_+^z e^{-iE_+ t/\hbar} + b \psi_-^z e^{-iE_- t/\hbar} \quad \psi(0) = \begin{pmatrix} a \\ b \end{pmatrix}$$

notation: $b := \sin(\alpha/2) \quad a := \cos(\alpha/2) \rightarrow |a|^2 + |b|^2 = 1$
normalized

$$\begin{aligned} \psi(t) &= \cos\left(\frac{\alpha}{2}\right) \psi_+^z e^{-iE_+ t/\hbar} + \sin\left(\frac{\alpha}{2}\right) \psi_-^z e^{-iE_- t/\hbar} = \\ &= \begin{pmatrix} \cos(\alpha/2) e^{-iE_+ t/\hbar} \\ \sin(\alpha/2) e^{-iE_- t/\hbar} \end{pmatrix} = \begin{pmatrix} \cos(\alpha/2) e^{i\gamma B t/2} \\ \sin(\alpha/2) e^{-i\gamma B t/2} \end{pmatrix} \end{aligned}$$

$$\begin{aligned} \langle S_x \rangle &= \psi^\dagger S_x \psi = \psi^\dagger(t) \cdot \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \cos(\alpha/2) e^{i\gamma B t/2} \\ \sin(\alpha/2) e^{-i\gamma B t/2} \end{pmatrix} = \\ &= \frac{\hbar}{2} \left(\cos\left(\frac{\alpha}{2}\right) e^{-i\gamma B t/2}, \sin\left(\frac{\alpha}{2}\right) e^{i\gamma B t/2} \right) \begin{pmatrix} \sin\left(\frac{\alpha}{2}\right) e^{-i\gamma B t/2} \\ \cos\left(\frac{\alpha}{2}\right) e^{i\gamma B t/2} \end{pmatrix} = \end{aligned}$$

$$= \frac{\hbar}{2} \cos\left(\frac{\alpha}{2}\right) \sin\left(\frac{\alpha}{2}\right) e^{-i\gamma B t} + \frac{\hbar}{2} \sin\left(\frac{\alpha}{2}\right) \cos\left(\frac{\alpha}{2}\right) e^{i\gamma B t/2} = \frac{\hbar}{2} \sin(\alpha) \cos(\gamma B t)$$

$2 \sin x \cdot \cos x = \sin(2x)$

$$= \frac{\hbar}{2} \sin(\alpha) \cos(\alpha/2) \left[\frac{e^{-i\gamma B t} + e^{i\gamma B t}}{2} \right] = \frac{\hbar}{2} \sin(\alpha) \cos(\gamma B t)$$

similarly

$$\langle S_y \rangle = -\frac{\hbar}{2} \sin(\alpha) \sin(\gamma B t)$$

using S_y and S_z

$$\langle S_z \rangle = \frac{\hbar}{2} \cos(\alpha)$$